

When you include time-dependence, quantum wave functions are solutions to the Schrödinger equation:

$$i\hbar \frac{\partial \Psi}{\partial t} = \hat{H} \Psi$$

The *Hamiltonian* operator \hat{H} is the total energy. With a 1D particle in a box, that's just the kinetic energy: $\hat{H} = \hat{K}$. All quantum states can be expressed as a superposition of the energy eigenfunctions we have found before. Say your wave function combines the first and second energy levels:

$$\Psi(x, t) = N [\psi_1(x)e^{-i\omega_1 t} + i\psi_2(x)e^{-i\omega_2 t}]$$

Here, ψ_n refers to the n th energy eigenfunction and $\omega_n = K_n/\hbar$. N is a normalization constant.

1. (30 points) Show that Ψ satisfies the Schrödinger equation. (Plug it in.)

Answer:

$$\begin{aligned} i\hbar \frac{\partial \Psi}{\partial t} &= iN\hbar \frac{\partial}{\partial t} [\psi_1(x)e^{-i\omega_1 t} + i\psi_2(x)e^{-i\omega_2 t}] = N\hbar [\omega_1 \psi_1(x)e^{-i\omega_1 t} + i\omega_2 \psi_2(x)e^{-i\omega_2 t}] \\ &= N [K_1 \psi_1(x)e^{-i\omega_1 t} + iK_2 \psi_2(x)e^{-i\omega_2 t}] \end{aligned}$$

Notice that $\frac{\partial^2}{\partial x^2} \psi_n = -(n\pi/L)^2 \psi_n = -2mK_n \psi_n/\hbar^2$. Therefore,

$$-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \Psi = N [K_1 \psi_1(x)e^{-i\omega_1 t} + iK_2 \psi_2(x)e^{-i\omega_2 t}]$$

These are the same, so Schrödinger's equation is satisfied.

2. (20 points) You know that the ψ_n are normalized:

$$\int_0^L dx \psi_1^* \psi_1 = \int_0^L dx \psi_2^* \psi_2 = 1$$

Show that ψ_1 and ψ_2 are *orthogonal*:

$$\int_0^L dx \psi_1^* \psi_2 = 0$$

Answer:

$$\int_0^L dx \psi_1^* \psi_2 = \frac{2}{L} \int_0^L dx \sin\left(\frac{\pi}{L}x\right) \sin\left(\frac{2\pi}{L}x\right) = 0$$

3. (50 points) Find N , the normalization constant.

Answer: Normalization means that

$$\int_0^L dx \Psi^* \Psi = N^* N \int_0^L dx [\psi_1^*(x)e^{i\omega_1 t} - i\psi_2^*(x)e^{i\omega_2 t}] [\psi_1(x)e^{-i\omega_1 t} + i\psi_2(x)e^{-i\omega_2 t}] = 1$$

There are four integrals to do. We can use normalization or orthogonality for all.

$$\begin{aligned} \int_0^L dx \psi_1^* \psi_2 e^{i\omega_1 t} e^{-i\omega_1 t} &= \int_0^L dx \psi_1^* \psi_2 = 1 \\ i \int_0^L dx \psi_1^* e^{i\omega_1 t} \psi_2 e^{-i\omega_2 t} &= i e^{-i(\omega_1 - \omega_2)t} \int_0^L dx \psi_1^* \psi_2 = 0 \\ -i \int_0^L dx \psi_2^* e^{i\omega_2 t} \psi_1 e^{-i\omega_1 t} &= i e^{-i(\omega_2 - \omega_1)t} \int_0^L dx \psi_2^* \psi_1 = 0 \\ -i^2 \int_0^L dx \psi_2^* \psi_2 e^{i\omega_2 t} e^{-i\omega_2 t} &= \int_0^L dx \psi_2^* \psi_2 = 1 \end{aligned}$$

Therefore

$$\int_0^L dx \Psi^* \Psi = N^* N (1 + 1) = 1$$

and

$$N = \frac{1}{\sqrt{2}}$$